

Spectral Theory.

i.e. $\partial_{xx} v + \partial_{yy} v = 0$ for $y > 0$.

Example. Dirichlet problem in Ω_t :

$$\left\{ \begin{array}{l} (\partial_{xx} + \partial_{yy}) v(x,y) = 0 \text{ in } \Omega_t = \{(x,y) : y > 0\} \\ v(x+i,0) = f(x) \text{ on } \mathbb{R} \\ v \text{ "not too large at } y = +\infty" \end{array} \right.$$

Formal solution. Set $v =$

$$\text{Observe that } -\partial_{xx} = A \geq 0$$

in the sense that for $f \in C_c^2(\mathbb{R}, \mathbb{C})$ with compact support we have

$$(Af, f)_{L^2} = \int_{-\infty}^{+\infty} -\partial_{xx} f \cdot \bar{f} dy =$$

$$= \int_{-\infty}^{+\infty} \partial_x f \cdot \partial_x \bar{f} dy \geq 0$$

If we could find A as a real number (or as a nonnegative function):

$$V(x,y) = \left(e^{-y\sqrt{A}} f \right)(x)$$

Then $V(x,0) = f(x)$

$$\begin{aligned} \partial_y V(x,y) &= \partial_y \left(\int_0^y e^{-t\sqrt{A}} f(t) dt \right) = A e^{-y\sqrt{A}} f \\ &= -\partial_{xx} V(x,y), \end{aligned}$$

OBS. We can make sense of this using Fourier transforms:

$$(Af)^*(t) = \int_{-\infty}^{+\infty} Af(x) e^{-2\pi ixt} dx$$

$$= (-2\pi i t)^2 \hat{f}(t) = 4\pi^2 t^2 \hat{f}(t)$$

and this suggests to consider

$$\begin{aligned} \hat{V}(t,y) &= e^{-2\pi i t y} \hat{f}(t) \text{ so that} \\ \partial_y \hat{V}(t,y) &= 4\pi^2 t^2 \hat{f}(t) e^{-2\pi i t y} \end{aligned}$$

$$\left. \begin{array}{l} \text{F. T. w.r.t.} \\ \text{then is} \\ \text{variables} \end{array} \right\} \text{and, after posing}$$

$$V(x,y) = \int_{-\infty}^{+\infty} \hat{V}(t,y) e^{2\pi i t x} dt \quad (\text{Fourier inversion})$$

$$\text{one has } \partial_y V(x,y) = A V(x,y) = -\partial_x V(x,y).$$

$$\text{Also: } \int_{-\infty}^{+\infty} |V(x,y)|^2 dx = \int_{-\infty}^{+\infty} |\hat{V}(t,y)|^2 dt =$$

$$= \int_{-\infty}^{+\infty} e^{-4\pi|t|y} |\hat{f}(t)|^2 dt \xrightarrow[y \rightarrow +\infty]{} 0$$

by M.C.T. if (e.g.) \hat{f} (linear, f) belongs to $L^2(\mathbb{R})$.

After we have formally solved the Dirichlet problem, we might want a more direct formula:

$$v(x, y) = \frac{1}{\pi} \int_{-\infty}^{+\infty} f(x-s) \frac{s}{s^2 + y^2} ds, \text{ a formula we have already used.}$$

Check it (exercise)

Hint. (i) $(f * g)^r = f^r \cdot g^r$

(ii) $\int_{-\infty}^{+\infty} e^{-2\pi t|y|} e^{2\pi i t x} dt =$

$$= \left[\frac{e^{2\pi t(ix-y)}}{2\pi i ix - y} \right]_0^\infty + \left[\frac{e^{2\pi t(ix+y)}}{2\pi i(ix+y)} \right]_\infty^0 \\ = \frac{1}{2\pi} \left(\frac{1}{ix+y} - \frac{1}{ix-y} \right) = \frac{2y}{\pi(x^2+y^2)}.$$

The advantage of the formal solution is that it works whenever $(Af, f) \geq 0$

If C^{∞} suitable Hilbert space.

e.g. $A = -\Delta_H$ where Δ_H is the Laplace-Beltrami operator on the Riemann manifold, etc.

Exercise.

Formally solve the

$$\begin{cases} \partial_y v(x, y) = \frac{\partial}{\partial x} v(x, y) & \text{for } x \in \mathbb{R}, y > 0 \\ v(x, 0) = \varphi(x) & \varphi \in L^2(\mathbb{R}) \\ \partial_y v(x, 0) = \psi(x) & \psi \in L^2(\mathbb{R}), \end{cases}$$

Observe that the solution works for $x \in H$, a Banach instead of a Hilbert.

Above we "computed":

$$\sqrt{-\Delta} \quad (\Delta = \frac{\partial^2}{\partial x^2})$$

$$e^{-yA} \quad (A = \sqrt{-\Delta})$$

$$\cos(yA), \sin(yA) \quad (\text{im the exercise}).$$

How do we make sense of them?

General idea: 1st compute a formal solution to the problem:
 2nd: make it less formal;
 3rd: study its properties.

Spectral Theory of Normal
and self-adjoint metrices.

$A \in M_n(\mathbb{C})$ is self-adjoint if $A = A^*$

it is unitary if $A^*A = I$

it is normal if $A^*A = A^*A$

unitary

self-adjoint \Rightarrow normal

then $A^* \in M_n(\mathbb{C})$: $(A^*v, v) := (v, Av)$

$\forall v \in \mathbb{C}^n$

$$D = \begin{pmatrix} d_1 & & 0 \\ 0 & d_2 & \\ & \ddots & \ddots \end{pmatrix} = D(d_1, \dots, d_n)$$

The goal: $D(d_1, \dots, d_n)$ is normal

and it is self-adjoint $\Leftrightarrow d_1, \dots, d_n \in \mathbb{R}$

if it is unitary ($\Rightarrow |d_1| = \dots = |d_n| = 1$)

Prop. $A \in M_n(\mathbb{C})$ is unitary

$\Leftrightarrow (Av, Aw) = (v, w) \quad \forall v, w \in \mathbb{C}^n$

$\Leftrightarrow |Av|^2 = |v|^2 \quad \forall v \in \mathbb{C}^n$

P.S. Unitary $\Leftrightarrow AA^* = I \Leftrightarrow$

$\forall v, w \in \mathbb{C}^n$: $(Av, w) = (A^*Av, w) = (Av, Aw)$

$$\Rightarrow |Av|^2 = |Av|^2 \quad \forall v \in \mathbb{C}^n$$

Polarization
Identity
Exercise

$$= |Av + Aw|^2 - |Av - Aw|^2 + i|Av + iw|^2 - i|Av - iw|^2$$

$$= 8(Av, Aw), \text{ by the same polarization.}$$

The characteristic polynomial
of $A \in M_{n \times n}(\mathbb{C})$ is $p(z) = \det(zI - A)$:

$$\deg(p) = n, \text{ hence, } p \text{ has } n$$

solutions $z = d_1, \dots, d_n$ (counting
multiplicities).

$$\text{Obs. } z = d_i \quad (\exists i) \Leftrightarrow zI - A \text{ is}$$

not invertible $\Leftrightarrow \exists v \in \mathbb{C}^n, v \neq 0$:

$Az = zv$ (z is an eigenvalue
of A and v is one of the
corresponding eigenvectors).

Theorem (I) $A \in M_n(\mathbb{C})$ is normal

$\Leftrightarrow \exists V$ unitary s.t. $A = V^{-1} D V$
for a diagonal matrix.

Moreover, if $D = D(d_1, \dots, d_n)$

then $\lambda_1, \dots, \lambda_n$ are the eigenvalues of A , and $\dim\{\text{vec}^n : Av = \lambda_j v\} = \mu_{\lambda_j}$

The dimension of the eigenspace relative to λ_j is the multiplicity of λ_j as root of $P(z) = 0$.

Moreover, $H(\lambda_j) \perp H(\lambda_i)$ if $\lambda_j \neq \lambda_i$.

(\square) If A is self-adjoint, the above holds and, moreover,

$\lambda_1, \dots, \lambda_n \in \mathbb{R}$. Viceversa,

if A is normal and $\lambda_1, \dots, \lambda_n \in \mathbb{C}$, then A is self-adjoint.

Sketch of the proof: (I) $A = V^* D V \Rightarrow$

$$A^* A = V^* \overline{D} V V^* D V = V^* D V^* D = V^* I D V =$$

$$= A A^*$$

Lemma: $A \in M_n(\mathbb{C}) \Rightarrow \exists V: V^* V = I$

and V^{-1} upper triangular.

$$V^{-1} =$$

$$\begin{bmatrix} t_{11} & & & \\ t_{12} & t_{22} & & \\ \vdots & \vdots & \ddots & \\ 0 & 0 & \cdots & t_{nn} \end{bmatrix}$$

$$\text{s.t. } A = V^* T V$$

from the Lemma: A is normal \Rightarrow

$$A^* A = V^* T^* T V$$

$$A A^* = V^* T T^* V \Rightarrow V^{-1} T^* = T V \Rightarrow T \text{ is diagonal}$$

Diagonalizability

$$\text{In fact } T^* T = \begin{bmatrix} t_{11} & 0 & 0 & \cdots & 0 \\ 0 & t_{22} & 0 & \cdots & 0 \\ 0 & 0 & t_{33} & \cdots & \vdots \\ \vdots & \vdots & \vdots & \ddots & \vdots \end{bmatrix}$$

$$= \begin{bmatrix} |t_{11}|^2 & & & & \\ & |t_{12}|^2 + |t_{22}|^2 & & & \\ & & |t_{13}|^2 + |t_{23}|^2 + |t_{33}|^2 & & \\ & & & \ddots & \vdots \end{bmatrix}$$

$$\text{while } T T^* = \begin{bmatrix} |t_{11}|^2 & & & & \\ & |t_{21}|^2 & & & \\ & & |t_{31}|^2 & & \\ & & & \ddots & \vdots \end{bmatrix}$$

Comparing elements on the diagonal:

$$t_{12} = \dots = t_{1n} = 0$$

$$t_{23} = \dots = t_{2n} = 0$$

i.e. T is diagonal.

Pf. of the Lemma. Let $v \in \mathbb{C}^n$, $|v| = 1$, s.t. $A v = \lambda v$ ($\exists \lambda$) and find

$$v = (v_1 v_2 \dots v_n) \text{ unitary, so that } (v_i v_j) = 0.$$

$$\text{Then } v^* A v = \left(\frac{\bar{v}_1}{\bar{v}_2}\right) A \left(v_1 v_2 \dots\right) =$$

$$= \left(\frac{\bar{v}_1}{\bar{v}_2}\right) (A v_1 A v_2 \dots) = \left(\frac{\bar{v}_1}{\bar{v}_2}\right) (\lambda v_1 \lambda v_2 \dots) =$$

$$= \left(\frac{\bar{v}_1}{\bar{v}_2}\right) (\lambda v_1 v_2 \dots) = \left(\frac{\bar{v}_1}{\bar{v}_2}\right) = \text{by induction} =$$

$$= \left(\frac{\bar{v}_1}{\bar{v}_2}\right) \left(\frac{\bar{v}_2}{\bar{v}_3}\right) \dots = \left(\frac{\bar{v}_1}{\bar{v}_n}\right) = \left(\frac{\bar{v}_1}{\bar{v}_2}\right) \left(\frac{\bar{v}_2}{\bar{v}_3}\right) \dots \text{ upper tr.}$$

$$= \left(\frac{1}{\partial t V^*} \right) \left(\frac{\lambda}{\partial t \bar{V}} \right) \left(\frac{1}{\partial t V} \right)$$

and the lemma is proved with
since $A = \left[\left(\frac{1}{\partial t V} \right) V^* \right]^k \left(\frac{\lambda}{\partial t \bar{V}} \right) \left[\left(\frac{1}{\partial t V} \right) V^* \right]^k$

$$\text{we have } \Pi_i \Pi_j = \begin{cases} \Pi_i & i=j \\ 0 & i \neq j \end{cases}$$

Since $A = V^* D(\lambda_1, \dots, \lambda_n) V$, the properties of eigenvectors/eigenspaces can be proved directly from $D(\lambda_1, \dots, \lambda_n)$. That's easy exercise. ■

Let $\lambda_1 = \dots = \lambda_{m_1} = \mu_1, \lambda_{m_1+1} = \lambda_{m_2} = \mu_2, \dots, \lambda_{m_l} = \mu_l$
be the distinct eigenvalues and π_1, \dots, π_l be the eigenspaces;

$\mathcal{C}^n \xrightarrow{\pi_j} H_j$ the orthogonal projection.

$$\Pi_j = V^* \begin{pmatrix} 0 & 0 & 0 \\ 0 & I_{m_j} & 0 \\ 0 & 0 & 0 \end{pmatrix} V$$
 where

$$I = \begin{pmatrix} I_m & & & \\ 0 & I_{m_2} & & \\ \vdots & \vdots & \ddots & \\ 0 & 0 & \cdots & I_{m_l} \end{pmatrix}$$

We can write $A = \sum_{j=1}^l \lambda_j \Pi_j$.

$\sigma(A) = \{\mu_1, \dots, \mu_l\}$ is the spectrum of A while $A = \sum_{j=1}^l \lambda_j \Pi_j$ is the spectral decomposition of A

$$\text{we have } \Pi_i \Pi_j = \begin{cases} \Pi_i & i=j \\ 0 & i \neq j \end{cases}$$

If φ is a polynomial,

$$\varphi(z) = \sum_{j=0}^n c_j z^j$$

$$\text{then } \varphi(A) = \sum_{j=0}^n c_j \left(\sum_{i=1}^l \lambda_i \Pi_i \right)^j =$$

$$= \sum_{j=0}^n c_j \sum_{i=1}^l \lambda_i^j \Pi_i = \\ = \sum_{i=1}^l \left(\sum_{j=0}^n c_j \lambda_i^j \right) \Pi_i = \sum_{i=1}^l \varphi(\lambda_i) \Pi_i$$

i.e. $\varphi(A)$ has the same spectral decomposition as A , with $\varphi(\lambda_i)$ instead of λ_i :

$$\sigma(\varphi(A)) = \varphi(\sigma(A)) \quad \begin{matrix} \text{Spectral} \\ \text{Mapping} \\ \text{Theorem} \end{matrix}$$

Also observe that

$$\varphi_1 = \varphi_2 \text{ on } \sigma(A) \Rightarrow \varphi_2(A) = \varphi_1(A)$$

i.e. operators of the form ' $\varphi(A)$ ' might be identified with ' $\sigma(A) \xrightarrow{\varphi} \varphi$ '.

Since $\forall \sigma(A) \xrightarrow{\phi} \sigma$ then is a polynomial q s.t. $q|_{\sigma(A)} = \phi$,

the spectral theory of normal matrices is simpler than the general one. (There are non-trivial problems).

$$\text{e.g. } A = \begin{pmatrix} 1 & 0 \\ 0 & e^{is} \end{pmatrix} = \pi_1 + e^{is}\pi_2$$

$$e^A = \sum_{n=0}^{\infty} \frac{(A_1 + e^{is}A_2)^n}{n!} = e^{\pi_1} + e^{e^{is}\pi_2} =$$

$$= q(A) \quad \text{where } q(1) = e^{\pi_1}$$

$$q(e^{is}) = e^{e^{is}\pi_2}$$

clearly we can realize it with $q(B) = \alpha x + b$ with $\alpha, b \in \mathbb{C}$.

In the infinite dimensional case things are not so straightforward, but most of the theory can be recovered, with useful applications.

§1 - Bounded Operators and adjoint.
 If Hilbert, $H \xrightarrow{T} H$ (inner) bounded
 Then $H \xrightarrow{\pi^*} H$, the adjoint of π^* ,
 is defined as: $(\pi^*x, y) := (x, \pi^*y)$.

$$(\pi^{-1})^* = (\pi^*)^{-1}; \quad \pi^{**} = \pi$$

$$\|\pi^*\pi\| = \|\pi\|^2; \quad \|\pi^*\pi\| \geq \sup_{x \neq 0} (\langle x, \pi^*\pi x \rangle)^{1/2} = \|\pi\|^2$$

$$\text{self-adj: } \pi = \pi^*$$

$$\text{orthogonal projection: } P = P^*, \quad P = P^*$$

$$\text{Normal: } \pi^*\pi = \pi\pi^*$$

Resolvent. $\Delta \in \rho(\pi) \Leftrightarrow \lambda I - \pi$ is $\neq 0$ \Rightarrow bijection, (b) with \Rightarrow biadj inv. \hookrightarrow bounded inverse.

$\Delta \in \sigma(\pi)$ (spectrum) iff $\Delta \notin \rho(\pi)$

For $\lambda \in \sigma(\pi)$ there are several possibilities:

- (i) $\exists x \in H, x \neq 0 : \pi x = \lambda x : \text{no HS - } \pi$
 is not injective. We say
- (ii) π : point spectrum
 and λ is an eigenvalue,
 x is a vector.

(ii) λ is not an eigenvalue and

Ran $\rho(\lambda T - \tau)$ is not dense.

(iii) λ is an eigenvalue and

residual Spectrum

Ran $\rho(\lambda T - \tau)$ is not H is dense.

Ran $\rho(\lambda T - \tau)$ is H is dense.

Spectral radius.

$$r(T) := \sup_{\lambda \in \sigma(T)} \|\lambda\|$$

$\lambda \in \sigma(T)$

Thm. $\lim_{n \rightarrow \infty} \|T^n\|^{\frac{1}{n}} = r(T)$

If $r(T) = \|T\|$: $r(T^n) = \|T^n\|$

Thm. (Phillip) $\sigma(T^*) = \overline{\sigma(T)}$

$$R_\lambda(T^*) = [R_\lambda(T)]^{*}$$

Theorem : $T^* = \|T\| \Rightarrow T^*$ has empty residual spectrum.

Spectral radius.

Theorem $\sigma(T) \neq \emptyset$.

$$\text{Pf. } \frac{1}{\lambda - T} = \frac{1}{\lambda} \frac{1}{1 - \frac{T}{\lambda}} = \frac{1}{\lambda} \frac{1}{1 - \frac{T}{\lambda}} = \frac{1}{\lambda} \left(\frac{1}{1 - \frac{T}{\lambda}} \right)^n$$

If we could do it, if $\forall \lambda$, there
 $\lambda \rightarrow \frac{1}{\lambda - T}$ would be "enveloping"
and bounded,

wince constant. ■

Properties of the Invariant
Subspaces.

§ 2. Spectral Thm.

$\{\pi_\alpha\}_{\alpha \text{ band in } H}$: projections on H
 is a projection valued measure

(p.v.m.) if $\overline{\pi}f = 0$; $\pi_{(-\epsilon, \epsilon)} = \pi_R f \pi_R$

$$\sigma = \bigcup_{n=1}^{\infty} \sigma_n (\text{disj}) \Rightarrow$$

$$\overline{\pi}f = \lim_{n \rightarrow \infty} \bigoplus_{j=1}^n \pi_{\alpha_j}$$

$$\pi_{\alpha_1}, \pi_{\alpha_2} = \pi_{\alpha_1 \cap \alpha_2}$$

Obs. $x, y \in H$: $\omega \mapsto (\pi_{\alpha(\omega)}, y)$

is a $(\sigma\text{-val.})$ meas.

Sup. $f: \mathbb{R} \rightarrow \mathcal{A}$ is bounded

Then $\exists! B$ s.t.

$$(Bx, y) = \int f(\lambda) d(\pi_\lambda x, y)$$

Specified Thus:

$A = A^*$ bounded $\rightarrow \{\pi_\alpha\}$ p.v.m.

$$\pi_\alpha = \chi_{\{\alpha\}}^A \quad \text{s.t.} \quad A = \int \lambda d\pi_\lambda$$

$$\text{Moreover: } f(A) = \int f(\lambda) d\pi_\lambda = \hat{\phi}(f)$$

Defines a map $\hat{f}: \mathbb{R} \rightarrow \hat{\phi}(f)$ Basis on H

(a) \hat{f} is a σ -hom.

(b) $\|\hat{f}(f)\|_{B(H)} \leq \|f\|_{L^\infty}$

$$(c) \hat{f}^*(x \mapsto x) = A$$

(d) $f_n \rightarrow f$ pointwise and

$$f_n f_m \in C \Rightarrow \hat{f}(f_m) \rightarrow \hat{f}(f_n)$$

in the strong topology

i.e. $\hat{f}(f_n)_h \rightarrow \hat{f}(f)_h$ in H .

$$(e) Ay = \lambda y \Rightarrow \hat{f}(f)y = f(\lambda)y$$

$$(f) f \geq 0 \Rightarrow \hat{f}(f) \geq 0, \text{ i.e.}$$

$$\langle \hat{f}(f)_h, h \rangle \geq 0 \text{ for all}$$

$$(g) BA = AB = \hat{f}(f)B = B\hat{f}(f)$$

Remarks on spectrum.

(1) When $\dim(H) = +\infty$, it can happen that $\det(\sigma(A), A = A^*$) but there is no $h \in H \setminus \{0\}$ s.t.

$$(A - \lambda I)h = 0$$

Note: it is possible that (H) has no solution $\lambda \in \sigma, h \in H \setminus \{0\}$. This is in sharp contrast with the case $\dim(H) < +\infty$.

Example: $H = L^2[0,1]$ and $A = M_x: f \mapsto f' \cdot x$

A is selfadjoint:

$$\int_{\mathbb{R}} \varphi(x) \psi(x) dx = \int_0^1 \varphi(x) (\psi'(x) - x\psi(x)) dx$$

hence $\sigma(A) \neq \emptyset$. But $A\varphi = \lambda\varphi$

$$h^2$$

means $x\varphi(x) = \lambda\varphi(x)$ d.e.x $\in [0,1]$

and this implies $\varphi = 0$ d.e.x $\in [0,1]$.

Exercise: show that $\sigma(M_x) = [0,1]$

(2) If $Ah = \lambda h$ has solutions

($A \in \mathcal{B}(H)$, $\det, h \in H, h \neq 0$)

then, if $M = \text{span}\{h\}$, we have that M is an invariant subspace of A :

$$(i) A(M) \subseteq M$$

$$(ii) 0 \notin M \subseteq H$$

$$(iii) M$$
 is closed

An other open problem in Operator Theory asks if all $T^* \in \mathcal{B}(H)$ ($\dim(H) = +\infty$) a unit invariant subspaces.

- If $\dim(H) < +\infty$ may the since that $(\lambda I - T^*)^{-1} = 0$ has solutions.

- If $A = A^*$ they do as a consequence of the spectral theorem (Exercise).