

Introduction

Canonical systems \leftrightarrow de Branges spaces

Louis de Branges (b. 1932)

1959-1965

de Branges, Hilbert Spaces of Entire Functions, 1968

R. Romanov, CS and de Branges spaces

Ch. Remling, Spectral theory of CS.

I. Spectral theory for CS \Leftarrow an overview.

What is a canonical system?

$$j Y'(x) = z H(x) Y(x), \quad x \in [0, L)$$

(L can be $+\infty$.)

- $Y = \begin{pmatrix} Y_+ \\ Y_- \end{pmatrix}$, $Y: [0, L) \rightarrow \mathbb{C}^2$ abs. continuous vector-valued function

- $H(x)$ - 2×2 matrix valued function

with real entries, $H \geq 0$, $H \neq 0$ on any set of positive measure, $H \in L^1[0, l] \forall l < L$.

- $j = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$ matrix imaginary unit, $j^2 = -I$

- $z \in \mathbb{C}$ - spectral parameter

The problem has unique solution for any initial data $Y(0) = Y_0 \in \mathbb{C}^2$.

$$Y = Y(x, z) = \begin{pmatrix} Y_+(x, z) \\ Y_-(x, z) \end{pmatrix}$$

Moreover, $\forall x \in (0, L)$ $Y_+(x, z)$, $Y_-(x, z)$ -
entire functions of z .

We can also consider this equation for
matrices: $Y M' = z H M$

The solution with $M(0, z) = I$ is called
the fundamental solution.

The matrix $M(L, z)$ (or $M(x, z)$) is called
the monodromy matrix.

$$M = \begin{pmatrix} \Theta_+ & \Phi_+ \\ \Theta_- & \Phi_- \end{pmatrix}, \text{ where } \Theta = \begin{pmatrix} \Theta_+ \\ \Theta_- \end{pmatrix}, \Phi = \begin{pmatrix} \Phi_+ \\ \Phi_- \end{pmatrix}$$

solve the equation $Y Y' = z H Y$ with $Y(0) = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$
and $Y(0) = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$.

Moreover, $\forall x \in (0, L)$ $Y_+(x, z)$, $Y_-(x, z)$ -
entire functions of z . (*) Π po maspyny monoprostanu
II. ST for CS - an q carried $H \in L^1[0, L]$
regular, Weyl's limit circle case

Two cases: $\rightarrow H \notin L^1[0, L]$ - singular, Weyl's
limit point case

Regular case: Let $Y(0) \in \mathbb{R}^2$

$E_x(z) := Y_+(x, z) + i Y_-(x, z)$, $x \in (0, L]$ -
entire function of the Hermite-Biehler class

$H \leftrightarrow E_L$ (or $\mathcal{H}(E_L)$ - dB space)

Inverse theorem: any E_L from some class corresp
to some ^(unique) Hamiltonian.

Moreover, $\mathcal{H}(E_x) \subset \mathcal{H}(E_L)$ isometrically, $0 < x < L$

Singular case: $H \leftrightarrow \mu$ - spectral measure
 μ - meas. on \mathbb{R} , $\int \frac{d\mu(t)}{t^2+1} < \infty$

$M = \begin{pmatrix} \Theta_+ & \Phi_+ \\ \Theta_- & \Phi_- \end{pmatrix}$. Then $\forall z \in \mathbb{C} \setminus \mathbb{R} \quad \forall x \in \mathbb{R}$

$$\exists q(z) = \lim_{x \uparrow L} \frac{z\Theta_+(x, z) + \Phi_+(x, z)}{z\Theta_-(x, z) + \Phi_-(x, z)}$$

q - Weyl coeff, independ. on ζ .

q is a Herglotz function: $\lim_{z \rightarrow \infty} q(z) = 0$, $z \in \mathbb{C}^+$
 $q(\bar{z}) = \overline{q(z)}$ (3)

Then $q(z) = pz + c + \int \left(\frac{1}{t-z} - \frac{t}{t^2+1} \right) d\mu(t)$

$p \geq 0, c \in \mathbb{R}, \int_{\mathbb{R}} \frac{d\mu(t)}{t^2+1} < \infty$

$H \rightsquigarrow q \rightsquigarrow \mu$. Inverses then says that any such μ corresponds to the unique Hamiltonian.

Example: $H(x) = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$

$y Y' = z H Y \quad \begin{cases} -Y_-' = z Y_+ \\ Y_+' = z Y_- \end{cases}$

Solution with $Y_+(0) = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad \Theta(x, z) = \begin{pmatrix} \cos xz \\ -\sin xz \end{pmatrix}$
 $Y_-(0) = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \quad \Phi(x, z) = \begin{pmatrix} \sin xz \\ \cos xz \end{pmatrix}$

$E_x(z) := \Theta_+(x, z) + i\Theta_-(x, z) = \cos xz - i\sin xz = e^{-ixz}$

Let $f \in L^2(\mathbb{C}^2, [0, L]) = L^2(H, [0, L])$

$f = \begin{pmatrix} f_+ \\ f_- \end{pmatrix} \xrightarrow{\mathcal{U}} \frac{1}{\sqrt{\pi}} \int_0^L (f_+(t) \cos zt - f_-(t) \sin zt) dt = \int_0^L f : [0, L) \rightarrow \mathbb{C}^2 : \int_0^L (H f, \delta t) dt$
 $\begin{matrix} \Theta_+(t, z) \\ -\Theta_-(t, z) \end{matrix}$

$= \frac{1}{2\sqrt{\pi}} \int_{-L}^L g(t) e^{-izt} dt \quad g \in L^2(-L, L)$
 $f_+(z) = \frac{g(z) + g(-z)}{2}$
 $f_-(z) = \frac{g(z) - g(-z)}{2}$

Thus, this mapping \mathcal{U} is a unitary operator from the space, associated to a Hamiltonian to some de Branges space (4)

$\mathcal{F} L^2(-L, L)$ - Paley-Wiener space PW_L -
 is a special case of a de Branges. In general
 we also construct a Fourier-type transform
 from $L^2(H, [0, L])$ to $\mathcal{H}(E_L)$
 (in the regular case)

II Examples of canonical systems.

1. Pólya equation
 Q 2×2 matrix function
 Q locally summable on $[0, L)$, $Q \in L^1_{loc}[0, L)$

$$JX' + QX = zX$$

Let X^0 be the solution of $J(X^0)' + QX^0 = 0$
 with $\det X^0 = 1$

$$X := X^0 Y \quad J(X^0)' Y + J X^0 Y' + Q X^0 Y = z X^0 Y$$

$$H = X_0^T X_0, \text{ rank} = 2 \quad A^T J A = J \text{ (even } \det A = 1)$$

2. Schrödinger equation

$$-y'' + qy = zy, \quad q \in L^1_{loc}[0, L)$$

Choose y_1, y_2 - solutions of $-y'' + qy = 0$

$$\text{with } y_1(0) = y_2'(0) = 1$$

$$y_1'(0) = y_2(0) = 0$$

$$\text{Then } y_1' y_2 - y_1 y_2' \equiv -1$$

$$Y := \begin{pmatrix} y_1 & y_2 \\ y_1' & y_2' \end{pmatrix}^{-1} \begin{pmatrix} y \\ y' \end{pmatrix} = \begin{pmatrix} y_2 y' - y_2' y \\ y_1' y - y_1 y' \end{pmatrix}$$

Then y is a solution of $-y'' + qy = zy$ iff

$$yY' = zHY, \quad H = \begin{pmatrix} y_1^2 & y_1 y_2 \\ y_1 y_2 & y_2^2 \end{pmatrix} \quad \text{rank} H = 1$$

$$(yY')_1 = y''y_1 - y_1''y = (qy - zy)y_1 - qy_1y = -zyy_1$$

$$y_1^2(y_2y' - y_2'y) + y_1y_2(y_1'y - y_1y_2') = -y_1^2y_2'y + y_1y_2y_1'y = y_1y_2[y_1'y_2 - y_1y_2'] = -1$$

3. String equation.

$$\rho(x) > 0, \quad \rho \in L_{loc}^1(0, L)$$

$$-y'' = z^2 \rho y, \quad Y := \begin{pmatrix} zy \\ y' \end{pmatrix}$$

$$yY' = z \begin{pmatrix} \rho & 0 \\ 0 & 1 \end{pmatrix} Y$$

IV. Some technicalities

Proposition: For any CS there exist an equivalent CS with $T_2 H = 1$.

$$\begin{cases} yY' = zHY, & Y(0) = Y_0 \\ yY' - \tilde{y}\tilde{Y}' = z \int_0^x H(t) Y(t) dt \end{cases}$$

Proof: $\tilde{y}(x) := \int_0^x t_2 H(t) dt$

The set $\{\tilde{\xi}(x) : \tilde{\xi}(x) = \tilde{\xi}(x') \text{ for some } x' \neq x\}$

$t_2 H(t) = 0 \Rightarrow H(t) = 0 \Rightarrow \tilde{\xi}$ increases strictly

$\tilde{H}(\tilde{\xi}(x)) := H(x)$ defines \tilde{H} on $(0, L')$

$L' = \tilde{\xi}(L)$ Let Y be a solution to

$$yY' = zHY$$

(6)

$\tilde{Y}(\xi(x)) := Y(x)$. Consider $\frac{\tilde{H}(t)}{\text{tr} \tilde{H}(t)}$ - new Hamiltonian with trace 1. $\tilde{H}_0(t) = \frac{\tilde{H}(t)}{\text{tr} \tilde{H}(t)}$

$$\begin{aligned} \int_0^s H_0(t) \tilde{Y}(t) dt &= \int_0^s \frac{\tilde{H}(t)}{\text{tr} \tilde{H}(t)} \tilde{Y}(t) dt = \\ &= \int_0^{\xi(s)} \frac{\tilde{H}(\xi(x))}{\text{tr} \tilde{H}(\xi(x))} \tilde{Y}(\xi(x)) \cdot \text{tr} \tilde{H}(x) dx = \\ &= \int_0^{\xi(s)} H(x) Y(x) dx = \frac{1}{z} \int_0^{\xi(s)} \mathcal{J} Y'(x) dx = \\ &= \frac{1}{z} (\mathcal{J} Y(\xi(s)) - \mathcal{J} Y(0)) = \frac{1}{z} (\mathcal{J} \tilde{Y}(s) - \mathcal{J} Y(0)), \end{aligned}$$

$$\mathcal{J} \tilde{Y} = z H_0 \tilde{Y}, \quad \text{tr} H_0 \equiv 1.$$

$H \rightsquigarrow H_0$ Regular case $\rightsquigarrow L' < \infty$
Singular case $\rightsquigarrow L' = +\infty$.

Th. Let H be a Hamiltonian with $\text{tr} H \equiv 1$.
Then the equation $\mathcal{J} Y' = z H Y$, $Y(0) = Y_0$,
has unique solution $Y = Y(x, z)$, for any
fixed $x \in [0, l]$ $Y_+(x, z)$, $Y_-(x, z)$ are entire
functions of z exponential type $\leq \infty$
and $\|Y(x, z)\| \leq \|Y_0\| e^{\alpha |z|}$.

(7)

Proof: $\mathcal{J}Y(x, z) = \mathcal{J}Y_0 + z \int_0^x H(t) Y(t, z) dt$

Usual iteration procedure

$$\mathcal{J}Y_1(x, z) := \mathcal{J}Y_0 + z \int_0^x H(t) Y_0 dt$$

$$\mathcal{J}Y_{n+1}(x, z) := \mathcal{J}Y_0 + z \int_0^x H(t) Y_n(t, z) dt$$

Then $\|\mathcal{J}Y_1(x, z) - \mathcal{J}Y_0\| \leq |z| \cdot \alpha \|Y_0\|$

Note that $\|H(t)\| \leq 1$

$$\|Y_{n+1}(x, z) - Y_n(x, z)\| \leq \frac{|z|^{n+1} \alpha^{n+1} \|Y_0\|}{(n+1)!}$$

$$Y_{n+1}(x, z) - Y_n(x, z) = z \int_0^x H(t) (Y_n(t, z) - Y_{n-1}(t, z)) dt$$

and induction.

Now, $Y_n(x, z) = Y_0 + \sum_{k=1}^n (Y_k(x, z) - Y_{k-1}(x, z))$

Y_n converges uniformly on $0 \leq x \leq L$
 $|z| < R$

to a solution. $Y_n(x, z)$ - pols in $z \Rightarrow$

$$\|Y(x, z)\| \leq \Rightarrow \text{solution is entire on } z$$

$$\leq \|Y_0\| + \sum_{k=1}^{\infty} \frac{|z|^k \alpha^k \|Y_0\|}{k!} = \|Y_0\| e^{\alpha|z|}$$

Rem. $Y_0 \in \mathbb{R}^2 \Rightarrow Y(x, z)$ is real for real z .